

**Hirlam/Aladin All Staff meeting
Reykjavík, 15–19 April 2013**

**Localized horizontal discretizations with
appropriate adjustment properties for the
ALADIN dynamics**

**Steven Caluwaerts, Daan Degrauwe, Piet Termonia, Pierre Bénard,
Fabrice Voitus and Jean-François Geleyn**

Motivation

Constraints

Formulation for
SWE

Consequences
of asymmetry

Conclusions

1. Motivation
2. Constraints
3. Formulation for Shallow Water Equations
4. Consequences of time-asymmetry
5. Conclusions

- Aladin/Arome/Harmonie is a semi-implicit, semi-Lagrangian *spectral* model
- When going to higher resolutions and larger domain sizes, we will face some *scientific* and *technical* challenges:
 - ◆ Representation of non-smooth fields (e.g. high-resolution orography) is problematic
 - ◆ The atmospheric reference state for the SI must be spatially homogeneous
 - ◆ Spectral transforms require domain-wide (MPI) communications

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 - ◆ The atmospheric reference state for the SI must be spatially homogeneous
 - ◆ Spectral transforms require domain-wide (MPI) communications
- Study the replacement of the spectral basis functions with local basis functions (finite elements)
- First focus on scientific impact
- This idea is not original (Staniforth, 1977)

- We want to keep as much as possible intact
 - ◆ necessary for a fair scientific comparison!
 - ◆ limited development cost

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- **Constraint 2: keep time step organization**

1. inverse FFT
2. physics
3. semi-Lagrangian interpolations
4. explicit dynamics
5. LBC treatment
6. forward FFT
7. solve Helmholtz equation in spectral space

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keep these!

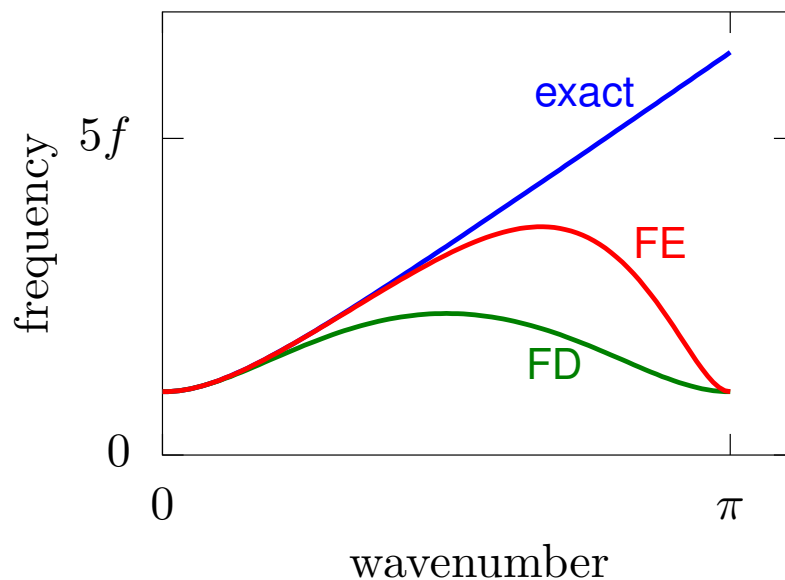
- Considering the linearized SWE equations in (u, v, h)

$$\frac{du}{dt} + g \frac{\partial h}{\partial x} + fv = 0$$

$$\frac{dv}{dt} + g \frac{\partial h}{\partial y} - fu = 0$$

$$\frac{dh}{dt} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0,$$

it is well known that localized schemes (finite differences/finite elements) on A-grid give *bad dispersion relations* for gravity wave propagation



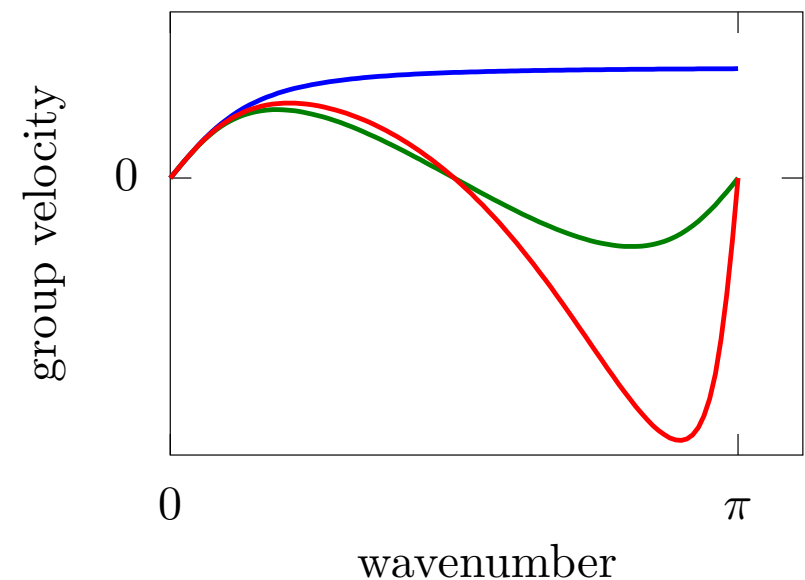
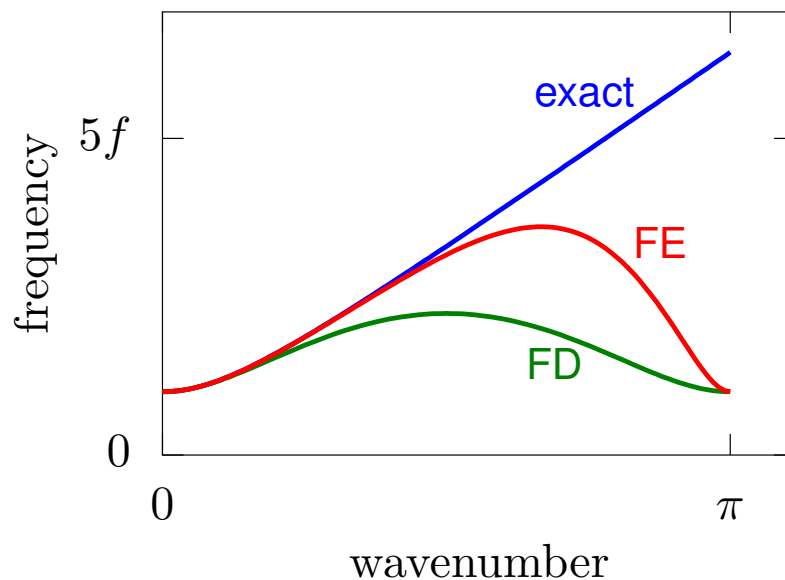
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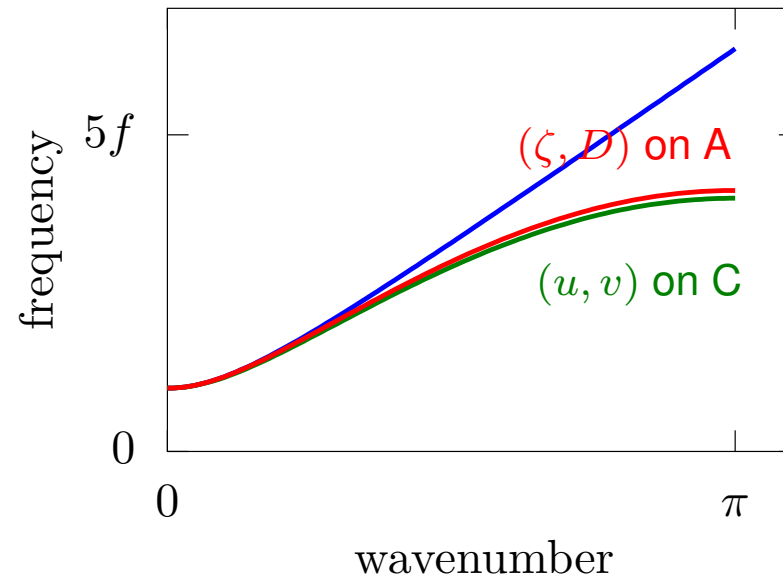
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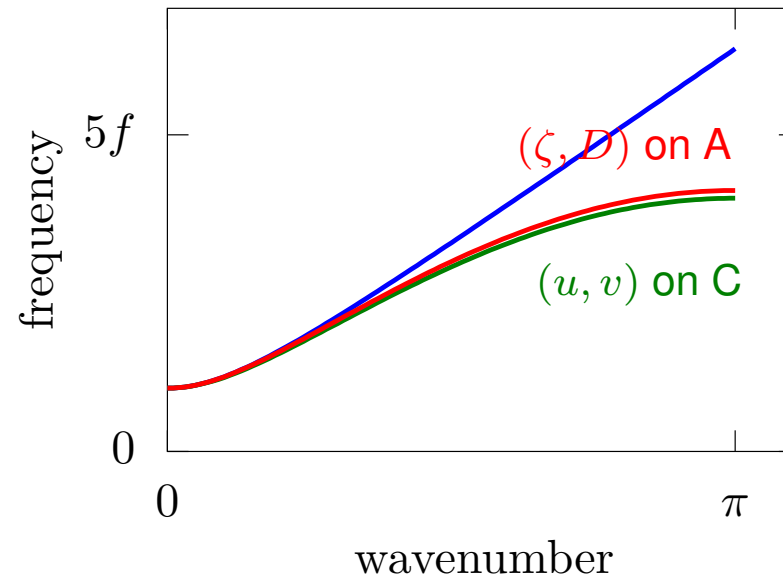


- Two possible solutions:



- ◆ Go to a staggered (C-) grid \Rightarrow not within constraints!
- ◆ Reformulate in terms of vorticity/divergence (ζ, D) instead of (u, v)
 \Rightarrow not entirely within constraints; we want to keep the RHS in (u, v)

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So we will try a *hybrid* $(u, v)/(\zeta, D)$ approach.

1. 3TL time discretization is done in (u, v)

$$u^+ + g\Delta t \left(\frac{\partial h}{\partial x} \right)^+ = R_u$$

$$v^+ + g\Delta t \left(\frac{\partial h}{\partial y} \right)^+ = R_v$$

$$h^+ + H\Delta t \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right)^+ = R_h$$

2. We switch to (ζ, D) without touching R_u , R_v and R_h

$$D^+ + g\Delta t \nabla^2 h^+ = \frac{\partial R_u}{\partial x} + \frac{\partial R_v}{\partial y}$$

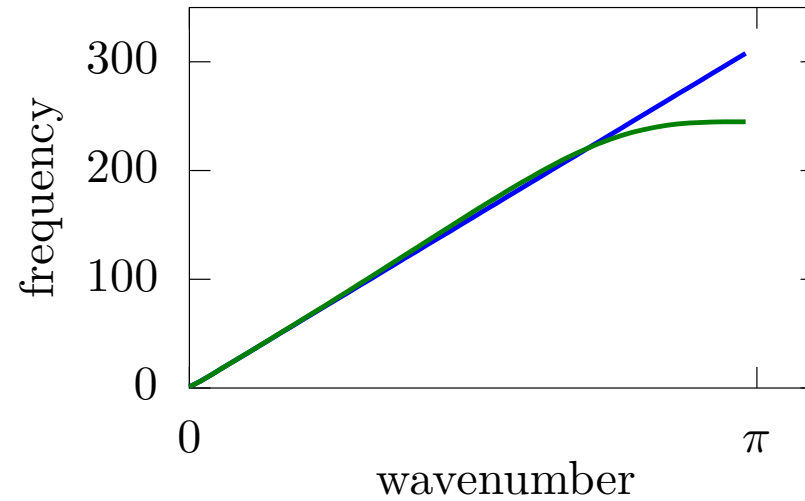
$$\zeta^+ = \frac{\partial R_v}{\partial x} - \frac{\partial R_u}{\partial y}$$

$$h^+ + H\Delta t D^+ = R_h$$

3. We solve this system with FE to (ζ^+, D^+, h^+)

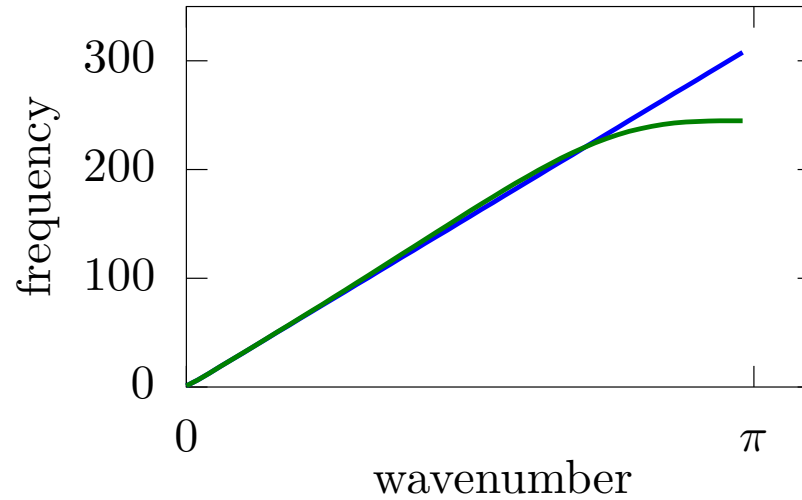
4. We transform (ζ, D) back to (u, v) , using FE

- The resulting dispersion relation looks okay

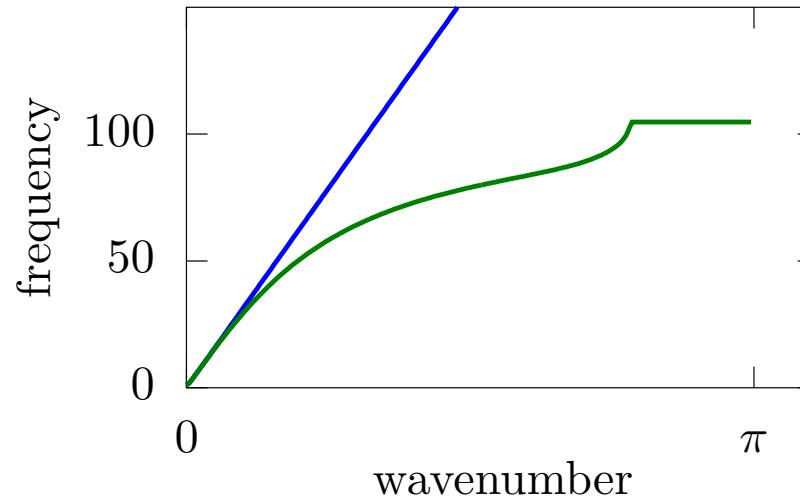


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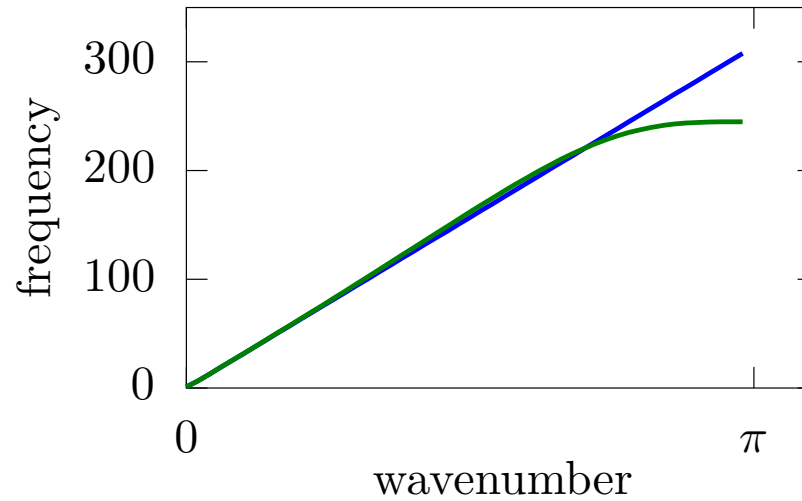
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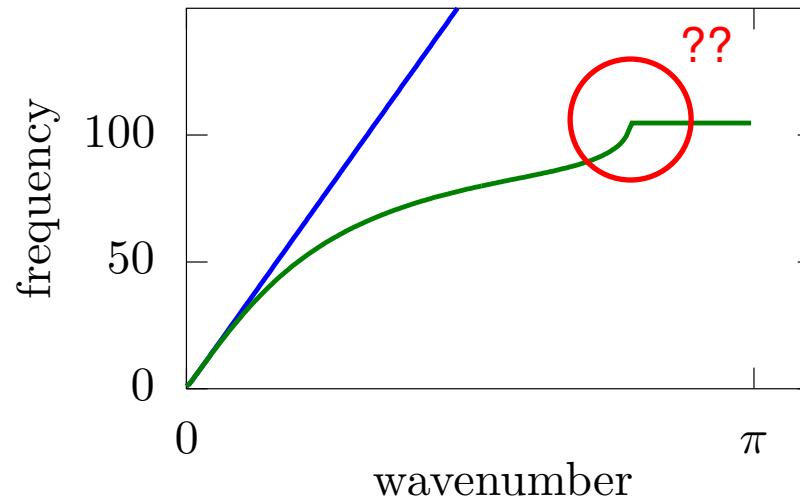
- But for larger timestep:



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- But for larger timestep:



- This behavior turns out to be the consequence of

1. with finite elements,

$$\frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial x} \right) \neq \frac{\partial^2 \psi}{\partial x^2}$$

2. there exists a time asymmetry in our scheme:

LHS contains $\nabla^2 h^+$

RHS contains $\frac{\partial R_u}{\partial x} + \frac{\partial R_v}{\partial y}$

with R_u containing $\frac{\partial h^-}{\partial x}$ and R_v containing $\frac{\partial h^-}{\partial y}$.

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- **Solution:** modify the pressure gradient terms in R_u and R_v to make the scheme symmetric again.
- **Remark:** such an asymmetry is also present in the Staniforth (1986) scheme, which performs poorly for *small* Δt .

- Resulting symmetrized hybrid $(u, v)/(\zeta, D)$ FE scheme:

$$\mathcal{S}_{xy}D^+ + g\Delta t(Q_x + Q_y)h^+ = \mathcal{L}_x\tilde{R}_u + \mathcal{L}_y\tilde{R}_v$$

$$\mathcal{S}_{xy}\zeta^+ = \mathcal{L}_x\tilde{R}_v - \mathcal{L}_y\tilde{R}_u$$

$$\mathcal{S}_{xy}h^+ + H\Delta t\mathcal{S}_{xy}D^+ = \mathcal{S}_{xy}R_h$$

$$\mathcal{S}_{xy}D^+ = \mathcal{L}_xu^+ + \mathcal{L}_yv^+$$

$$\mathcal{S}_{xy}\zeta^+ = \mathcal{L}_xv^+ - \mathcal{L}_yu^+$$

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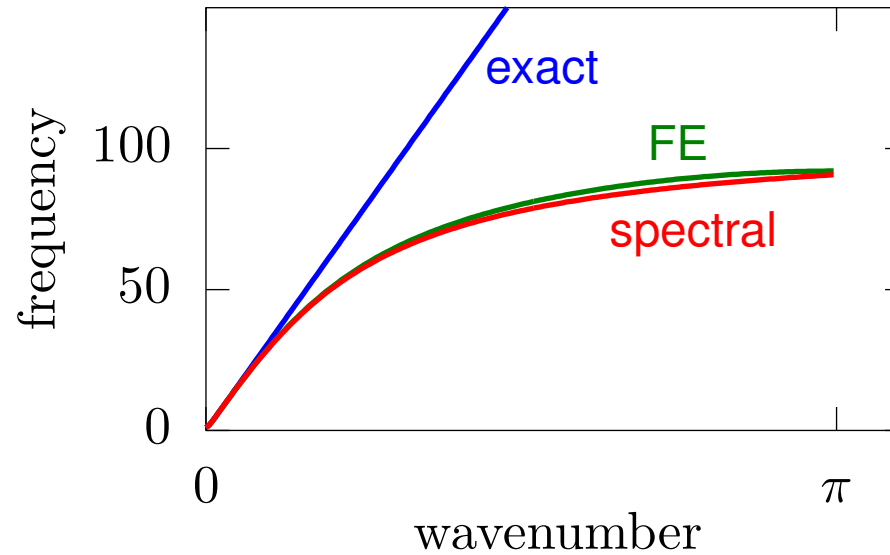
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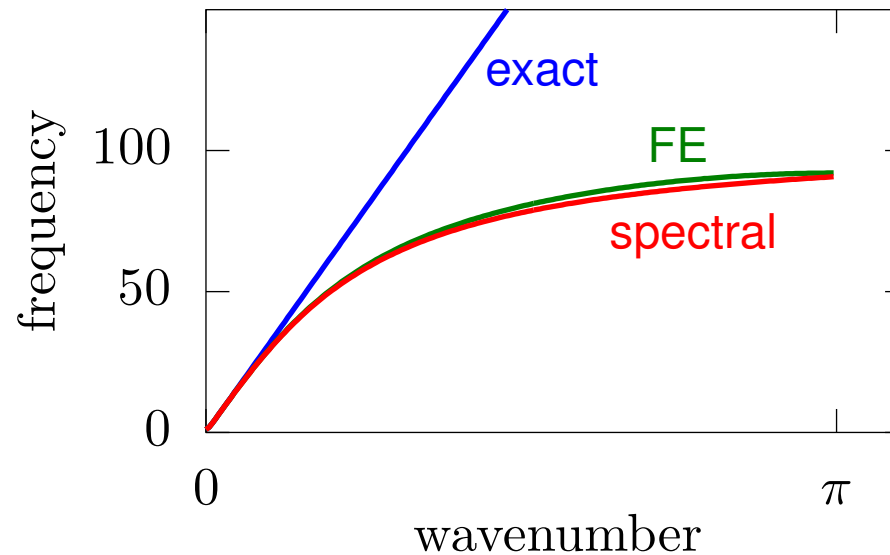
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- Dispersion relation:



- Note that a 2TL variant of the scheme can also be derived.

We present a localized horizontal discretization scheme:

- finite-element based
- on our non-staggered A-grid
- which fits into the ALADIN semi-implicit, semi-Lagrangian algorithmics
- doesn't require modifications to gridpoint calculations (physics!)
- has an excellent dispersion relation for gravity waves
- opens the way for answering scientific questions like a non-homogeneous reference state and influence of steep orography

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Thank you !