Regional Cooperation for Limited Area Modeling in Central Europe



ACC and RD A Consortium for COnvection-scale modelling Research and Development

# **Dynamics in LACE - towards hectometric scales**

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- Dynamical core in ACCORD
- **SI** time scheme
- Orographic terms in linear model (based on work of Jozef Vivoda and Fabrice Voitus)
  - **Some equations**
  - Idealised tests
  - Real simulations @200m
- □ Vertical velocity definition in nonlinear model (based on work of Fabrice Voitus)
  - Real simulations @200m







- □ hydrostatic primitive equation system (HPE) or Euler equations (EE); recently implemented quasi elastic equation system (QE)
- **D** prognostic variables  $\vec{v}, T, q_s = \ln(\pi_s)$ , in EE with  $w, \hat{q} = \ln(\frac{p}{\pi})$

Discretization

- **u** spectral method for horizontal direction
- □ hybrid vertical coordinate  $\eta$  based on hydrostatic pressure  $\pi(\eta) = A(\eta) + B(\eta)\pi_s$ ; A(top) = B(top) = 0, A(bottom) = 0, B(bottom) = 1
- □ finite differences or finite elements for vertical direction discretization
- □ semi-implicit or iterative centred implicit scheme for time discretization

semi-Lagrangian advection



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# Semi-Implicit time scheme



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System evolutionLinearization
$$\frac{dX}{dt} = \mathcal{M}X$$
 $X = X^* + X', \quad \mathcal{M} \to \mathcal{L}^*$ Using linear model  $\mathcal{L}^*$  we get $\frac{dX}{dt} = \mathcal{L}^*[\overline{X}]^t + (\mathcal{M} - \mathcal{L}^*)X$ Semi-implicit schemeand discretize in time to obtain $\frac{X^+ - X^0}{\Delta t} = \mathcal{L}^*\left(\frac{X^+ + X^0}{2}\right) + (\mathcal{M} - \mathcal{L}^*)X^{+\frac{1}{2}}$ Iterative centered implicit schemeor $\frac{X^{+(n)} - X^0}{\Delta t} = \frac{\mathcal{L}^*X^{+(n)} + \mathcal{L}^*X^0}{2} + \frac{(\mathcal{M} - \mathcal{L}^*)X^{+(n-1)} + (\mathcal{M} - \mathcal{L}^*)X^0}{2}$ We know that both can be second order accurate in time when some care is taken (averaging

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along semi-Lagrangian trajectory).



# **Full model**



#### Temperature

$$\frac{dT}{dt} = \frac{\kappa T}{\kappa - 1} \left( D + d \right)$$

## Horizontal momentum

$$\frac{d\vec{v}}{dt} = -RT\frac{\nabla\pi}{\pi} - \nabla\phi - RT\nabla\hat{q} - \frac{1}{m}\frac{\partial(p-\pi)}{\partial\eta}\nabla\phi$$

Vertical	momentum			
	$\frac{dw}{dt} =$	$\frac{g}{m}\frac{\partial(p)}{\partial q}$	$(D-\pi)\over\partial\eta$	

### Pressure departure

$$\frac{d\hat{q}}{dt} = \frac{1}{\kappa - 1} \left( D + d \right) - \frac{1}{\pi} \frac{d\pi}{dt}$$

## Surface pressure

$$rac{dq_s}{dt} \;=\; -rac{1}{\pi_s} \int_0^1 
abla \cdot (m ec v) d\eta$$

## Diagnostic relations

$$\frac{d\pi}{dt} = \vec{v} \cdot \nabla \pi - \int_{0}^{\eta} \nabla \cdot (m\vec{v}) d\eta'$$
$$\phi = \phi_{s} - \int_{\eta}^{1} \frac{mRT}{p} d\eta'$$
$$d = -\frac{p}{mRT} \left( \nabla \phi \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta} \right)$$

Definitions			
	D	=	$ abla \cdot ec v$
	$\kappa$	=	$\frac{c_p}{R}$
	m	=	$rac{\partial\pi}{\partial\eta}$

# **Full model**



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Definitions			
	D	=	$ abla \cdot ec v$
	$\kappa$	=	$\frac{c_p}{R}$
	m	=	$rac{\partial \pi}{\partial \eta}$

Current state:

stationary

resting

 $\Box$  hydrostatically balanced ( $\pi_s^*$ )

🖵 dry

 $\Box$  isothermal ( $T^*$ )

 $\Box$  with constant orography ( $\nabla \phi^* = 0$ )









# **Basic state**

New:

stationary

resting

 $\Box$  hydrostatically balanced ( $\pi_s^*$ )

🖵 dry

□ isothermal (*T*<sup>\*</sup>)

 $\Box$  with constant orographic slope (in absolute value,  $|\nabla \phi^*| \neq 0$ )









# Linear model



#### Temperature

$$\frac{\partial T}{\partial t} = \frac{\kappa T^*}{\kappa - 1} \left( D + d \right)$$

### Horizontal momentum

$$\frac{\partial \vec{v}}{\partial t} = -RT^* \frac{\nabla \pi}{\pi^*} - \nabla \phi - RT^* \nabla \hat{q} - \frac{1}{m^*} \frac{\partial \pi^* \hat{q}}{\partial \eta} \nabla \phi^*$$

Vertical	momer	itum
	$\frac{\partial w}{\partial t} =$	$= rac{g}{m^*} rac{\partial \pi^* \widehat{q}}{\partial \eta}$

#### Pressure departure

$$\frac{\partial \hat{q}}{\partial t} = \frac{1}{\kappa - 1} \left( D + d \right) + \frac{1}{\pi^*} \int_0^{\eta} m^* D d\eta'$$

### Surface pressure

$$rac{\partial q_s}{\partial t} = -rac{1}{\pi_s^*} \int_0^1 m^* D d\eta$$

#### Diagnostic relations

$$\nabla \phi = \nabla \phi_s - \int_{\eta}^{1} \nabla \left(\frac{mRT}{p}\right) d\eta'$$
  

$$\nabla \phi^* = g \Lambda^* S^* (\eta)$$
  

$$d = -\frac{p}{mRT} \left(\nabla \phi \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta}\right)$$

Definitions  $\Lambda^* = \frac{1}{g} ||\nabla \phi_s||^*$   $S^*(\eta) = \frac{B(\eta)\pi_s^*}{\pi^*(\eta)}$   $m^* = \frac{\partial \pi^*}{\partial \eta}$ 

# Linearized slope in SI



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Modified vertical divergence  $d = \frac{p}{mRT} \left( \nabla \phi \ \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta} \right)$ 

Time evolution in linear model

$$\begin{aligned} \frac{\partial \vec{v}}{\partial t} &= \mathbb{A} - \Lambda^* S^* (\eta) \mathbb{B} \\ \frac{\partial w}{\partial t} &= \mathbb{B} \\ \frac{\partial d}{\partial t} &= \frac{1}{RT^*} \left[ \nabla \phi^* \ \partial^* \left( \frac{\partial \vec{v}}{\partial t} \right) - g \partial^* \left( \frac{\partial w}{\partial t} \right) \right] \\ &= \frac{1}{RT^*} \left[ g \Lambda^* S^* (\eta) \partial^* \mathbb{A} - g \Lambda^{*2} S^* (\eta) \left( S^* \partial^* \mathbb{B} + \mathbb{B} \partial^* S^* \right) - g \partial^* \mathbb{B} \right] \end{aligned}$$
where
$$\partial^* X &= \frac{\pi^*}{m^*} \frac{\partial X}{\partial \eta}$$

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We omit the first order terms in  $\Lambda^*$  and then  $\frac{\partial \vec{v}}{\partial t}$  is unchanged and all operators of the RHS of  $\frac{\partial d}{\partial t}$  apply on  $\hat{q}$ .

Time evolution in linear model

$$\frac{\partial \vec{v}}{\partial t} = \mathbb{A} - \bigwedge^* S^*(\eta) \mathbb{B}$$

$$\frac{\partial d}{\partial t} = \frac{1}{RT^*} \left[ g \bigwedge^* S^*(\eta) \partial^* \mathbb{A} - g \bigwedge^{*2} S^*(\eta) \left( S^* \partial^* \mathbb{B} + \mathbb{B} \partial^* S^* \right) - g \partial^* \mathbb{B} \right]$$

Finally, since  $\mathbb{B} = g \left( \partial^* + 1 \right) \hat{q}$ 

Time evolution in linear model

$$\frac{\partial \vec{v}}{\partial t} = \mathbb{A} \quad \rightsquigarrow \quad \frac{\partial D}{\partial t} = \frac{\partial (\nabla \cdot \vec{v})}{\partial t} = \nabla \cdot \mathbb{A}$$
$$\frac{\partial d}{\partial t} = \mathcal{L}^*_{new} \hat{q}$$







We can define



How to discretize the proposed solution?

New discretized vertical Laplacian operator  $\begin{aligned} [\partial^*(\partial^* + 1)X]_l &= \dots \\ [(\partial^* + 1)X]_l &= \dots \\ S^*(\eta_l) &= \dots \\ \partial^*S^*(\eta_l) &= \dots \end{aligned}$ 

How to set boundary conditions?







Does  $\mathcal{L}_{new}^*$  have only real and negative eigenvalues?

For an example of 87 vertical levels used in Czech operations we are safe.



Then we can eliminate the discretized equations up to horizontal divergence D and solve the Helmholtz equation for D.





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Vertical velocity for the Schär mountain case depending on  $\Lambda^*$ . ( $\Delta x = 500 \, m, \Delta z = 250 \, m$ , mountain height  $h = 250 m, T_0 = 288K, u_0 = 10 m/s, \Delta t = 32 s$ )



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# **Real simulations @200m**





14/25 ASW 2024











The basic algorithmic choices for ALARO configurations @200m are:

### Dynamical core

- □ semi-Lagrangian advection scheme with 4 iterations for trajectory calculation
- □ PC time scheme with one iteration, cheap variant (SL trajectories are not recalculated in corrector)
- modified vertical divergence d4 for vertical motion, transformation to vertical velocity w in the non-linear model
- □ reference values of the linear model: SITR=300K, SITRA=100K, SIPR=900hPa
- **no** decentering
- □ semi-Lagrangian horizontal diffusion applied on all model variables + TKE, TTE, hydrometeors
- Inear truncation for all spectral fields except orography; quadratic truncation of orography





## ALARO physics

- □ radiation scheme ACRANEB2
- □ turbulence and shalow convection scheme TOUCANS, model 2
- □ scale aware deep convection and microphysics scheme 3MT

Initialization

□ initialization with 3DVAR + surface DA (canari) for 2.325km run; dynamical adaptation + DFI for 500m and 200m runs

Particular choices for ALARO@200m:

- cubic truncation of orography
- □ SITRA=50K
- □ no 3MT (deep convection), only STRAPRO (stratiform precipitation)







### Vertical velocity for the alpine case 19 August 2022 OUTC + 24hours.







Vertical velocity depending on  $\Lambda^*$  for the alpine case 19 August 2022 OUTC + 24hours.



With additional iterations of the SI scheme, the integration crashes.





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Vertical velocity depending on  $\Lambda^*$  for the alpine case 19 August 2022 OUTC + 24hours.



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# **Real simulations**



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Averaged spectral norms of vertical divergence for the alpine case 19 August 2022 OUTC + 24 hours.





Conclusions

- Use must continue our efforts.
- $\Box$  We plan to test various possible definitions of vertical function  $S^*(\eta)$ .
- □ We plan to test various possible discretizations including boundary conditions.
- U We plan to make further idealised tests and real simulations.
- □ If the time scheme allows a source of a noise, further time iterations may not help to stabilize the scheme. To the contrary, the scheme with further iterations may show even less stability.





# Vertical motion variable



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Fabrice Voitus proposed a modification of the vertical velocity defined in the model to simplify the bottom boundary condition and allow more precise calculation.

https://events.ecmwf.int/event/167/contributions/1379/attachments/794/1401/AS2020-Voitus.pdf

Current state	Fabrice's definition
$w = \frac{dz}{dt}$	$W \;=\; w - \mathbf{V} \cdot S(\eta)  rac{1}{g}  abla \phi_s$
$d = -g \frac{p}{mRT} \frac{\partial w}{\partial \eta} + X$	$d = -g \frac{p}{mRT} \frac{\partial W}{\partial \eta} + X$
$X = \frac{p}{mRT} \frac{\partial \mathbf{V}}{\partial \eta} \nabla \phi$	$X = \frac{p}{mRT} \frac{\partial \mathbf{V}}{\partial \eta} \nabla \left(\phi - S(\eta) \phi_s\right) - \frac{p}{mRT} \mathbf{V} \frac{\partial \nabla \phi}{\partial \eta}$
$\frac{\partial d}{\partial t} = -\frac{g^2}{RT^*} \mathcal{L}_v^* \hat{q}$	$\frac{\partial d}{\partial t} = -\frac{g^2}{RT^*} \mathcal{L}^*_{mod} \hat{q}$
$\mathcal{L}_v^* \ = \ \partial^* \left( \partial^* + 1  ight)$	$\mathcal{L}_{mod}^{*} \;=\; \partial^{*}\left(\partial^{*}+1+\Lambda^{*2}S\left(\eta ight)\gamma^{*} ight)$

Then discretization of these terms involves interpolations between half levels (where w is represented) and full levels (where V is represented) and is cumbersome ...

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#### Current state

in the non-linear model

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$$w_s = \mathbf{V}_s \frac{1}{g} \nabla \phi_s$$
  
 $\frac{dw_s}{dt} = \frac{d \left( \mathbf{V}_s \frac{1}{g} \nabla \phi_s \right)}{dt}$ 

while in the linear model explicit guess of d is calculated consistently and used in the implicit part

Complicated and not exact! It is a source of noise which may grow.

#### Fabrice's definition

in the linear and non-linear model

$$W_s = 0$$

$$\frac{dW_s}{dt} = 0$$

Easily applicable!

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#### Vertical velocity for the alpine case 19 August 2022 OUTC + 24hours.



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# **Real simulations**

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Averaged spectral norms of vertical divergence for the alpine case of 19 August 2022<sup>central europe</sup> OUTC + 24 hours.





Conclusions and advertisements

- □ The modification is available in cycle CY49t1 under namelist option NVDVAR=5, thanks to Fabrice Voitus and Karim Yessad.
- □ The new formulations may help to further reduce the non-linear residual of the ICI time scheme and to get rid of the noise coming from steep orography, especially in high resolutions.





# Tack för din uppmärksamhet!

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